

A Note on the Mathematical Model on a Cigarette-Like Combustion

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Abstract

This paper presents a mathematical model of a cigarette-like combustion using high activation energy asymptotic of particular interest are question of existence and uniqueness. It is shown that the steady temperature increase as the Frank-Kamenetskii λ_1 increases.

1.0 Introduction

Many scientists have contributed to the theory of combustion. Some have done a lot particularly on the criteria for existence and uniqueness of solution of some combustion problems [1-9].

Ayeni and Agosto [1] investigated the existence and uniqueness of a solution of self-similar diffusion equations. They examined the models of fast flow of a gas in a porous medium and microwaves heating of various materials. They gave the criteria for the existence and uniqueness of solution of the problem. Koriko and Ayeni [6] investigated the existence and unique of a system of partial differential equations arising from cigarette-like combustion.

Their results are useful in making deductions on porous medium combustion. Olanrewaju and Ayeni [8] investigated the effect of Frank-Kamenetskii parameter on strong detonations in converging vessels in which they established the fact that there is an appreciable difference in the temperature along the vessel when Frank-Kamenetskii parameters differ only by 1/30. In this paper, we extend the model of Koriko and Ayeni [6] to a situation when the Frank-Kamenetskii λ_1 and λ_2 are not equal.

2.0 Mathematical Formulation

A cigarette-like combustion is a two stage reaction [6] the two stages are first the diffusion of gas (air) between the gas mainstream and the reaction sites in the solid and secondly the conventional Arrhenius reaction.

$$\frac{\kappa a^2 T}{dx^2} + \frac{(T-T_0)^n Q A e^{-E/RT}}{a + b e^{-E/RT}} = 0 \tag{1}$$

As boundary conditions we take $T(-L) = T(L) = T_0$ (2)

Where:

- T is the temperature of the medium,
- T_0 is the initial temperature,
- K is the thermal conductivity,
- Q is the combustion energy,
- A is the pre-exponential factor,
- E_1 is the activation energy of the first stage ,
- E is the activation energy of the second reaction,
- a and b are permeability constants(dimensionless),
- n is the other of reaction.

We non-dimensionalize the temperature by writing:

$$T = \frac{E}{RT_0} (T - T_0), \quad y = \frac{x}{L} \text{ and we assumed } E = E_1$$

$$\frac{d^2 \theta}{dy^2} + \frac{\lambda_1 \theta^n e^{\frac{\theta}{1+\epsilon \theta}}}{1 + \lambda_2 \frac{\theta}{e^{\frac{\theta}{1+\epsilon \theta}}}} \tag{3}$$

$$\theta(-1) = \theta(1) = 0 \tag{4}$$

Where:

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$$\lambda_1 = \frac{L^2 \varepsilon^n Q A e^{-\frac{E}{RT_0}}}{a} \quad - \quad \text{Frank-Kamenetskill}$$

$$\lambda_2 = \frac{b}{a} e^{-\frac{E}{RT_0}} \quad - \quad \text{Permeability Parameter}$$

$$\varepsilon = \frac{RT_0}{E} \quad - \quad \text{Dimensionless Parameter}$$

$L = \text{half length of the medium}$

In the this paper, $n = 0, E = E_1$. The case $n \neq 0, E \neq E_1$ is the subject of another paper.

Method of Solution

We let:

$$\begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} y \\ \theta \\ \theta' \end{pmatrix} \tag{5}$$

Then we obtain

$$\begin{pmatrix} x_1' \\ x_2' \\ x_3' \end{pmatrix} = \begin{pmatrix} 1 \\ \theta' \\ \theta'' \end{pmatrix} \tag{6}$$

Where

$$\begin{pmatrix} x_1' \\ x_2' \\ x_3' \end{pmatrix} = \begin{pmatrix} 1 \\ x_3 \\ \frac{x_2}{-\lambda_1 e^{1+\varepsilon x_2}} \\ \frac{x_2}{1+\lambda_2 e^{1+\varepsilon x_2}} \end{pmatrix} \tag{7}$$

The initial conditions are:

$$\begin{pmatrix} x_1(-1) \\ x_2(-1) \\ x_3(-1) \end{pmatrix} = \begin{pmatrix} x_1(-1) \\ \theta(-1) \\ \theta(-1) \end{pmatrix} = \begin{pmatrix} -1 \\ 0 \\ \beta \end{pmatrix} \tag{8}$$

Let

$$Z = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} \tag{9}$$

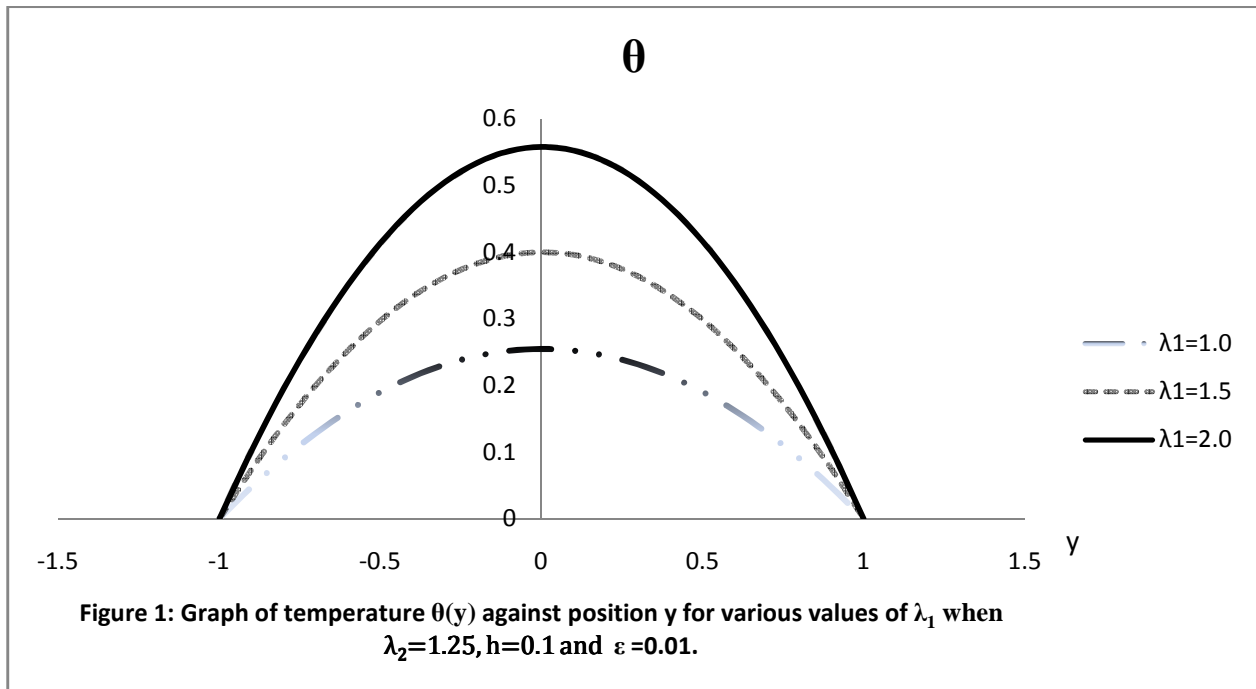
and

$$E(z) = \begin{pmatrix} 1 \\ x_3 \\ \frac{x_2}{-\lambda_1 e^{1+\varepsilon x_2}} \\ \frac{x_2}{1+\lambda_2 e^{1+\varepsilon x_2}} \end{pmatrix}$$

Then $z^1 = E(z)$ and

$$z(-1) = \begin{pmatrix} -1 \\ 0 \\ \beta \end{pmatrix}$$

We guess β . We continue to change β . The required β is the one for which $\theta(1) = 0$ (shooting method). We present in the figure 1 the graph of θ against y for various values of λ_1 .



3.0 Discussion of the Result

We observe that the maximum temperature increases with the increase in the frank-Kamenetskiis parameter λ_1 . Koriko [5] showed there are multiple solutions when only the initial and end points are fixed. The figure 1 shows that there exists only one solution when the derivative $\theta'(-1)$ is fixed. That is, the problem for which $\theta(-1) = 0$, $\theta'(-1) = \beta$ and $\theta(1) = 0$ is unique.

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