

## Projectile Motion in a Resistant Medium under the Influence of an Electric Field

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### Abstract.

The motion of a particle which is projected into a resistant medium and subjected to a uniform gravitational field and electric field is considered. The drag force that acts upon the particle within the medium is proportional to the speed raised to power  $n$ . The problem is formulated in terms of particle - speed and local-path-angle variables, and the equations of motion that result into non-linear and coupled. Of particular interest is the effect of the magnetic field on the velocity.

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### 1.0 Introduction

In a recent paper Hayen [2] investigated projectile motion in a resistant medium. The particle is subjected to a uniform gravitational field. The drag force that acts upon the particle within the medium is proportional to the particle speed squared. Although the physical problem considered in [2] has been previously solved, the particular solution technique presented and certain solution properties uncovered are original.

In another paper Vajravelu and Rivera [3] investigated the hydro-magnetic flow at an oscillating plate and at a moving plate. They discussed the effect of the magnetic parameter on the flow.

In this paper we extend the model of Hayen [1] to include an electric field and general power  $n$  which is new and we study the resulting unsteady problem.

### 2.0 Mathematical Formulation

The set of coupled non-linear equations are

$$\begin{aligned} \frac{dv}{dt} &= -g \sin \phi + \frac{\sigma}{\rho} B_0^2 v - \frac{k}{m} v^n, \quad v(0) = v_0 \\ v \frac{d\phi}{dt} &= -g \cos \phi, \quad \phi(0) = \phi_0 \end{aligned} \quad (2.1)$$

where  $v$  is the velocity,  $\phi$  the local path angle,  $g$  the acceleration due to gravity,  $m$  the mass of the particle and  $k$  is a constant. Also,  $\rho$  is the fluid density,  $\sigma$  the electric conductivity,  $B_0$  the imposed magnetic field,  $n$  a positive constant,  $x$  and  $t$  variables respectively.

### 3.0 Method of Solution

Consider the system of equation

$$\begin{aligned} x'_1 &= f_1(x_1, \dots, x_n, t), x_1(t_0) = x_{10} \\ x'_2 &= f_2(x_1, \dots, x_n, t), x_2(t_0) = x_{20} \\ &\vdots \\ x'_n &= f_n(x_1, \dots, x_n, t), x_n(t_0) = x_{n0} \end{aligned} \tag{3.1}$$

Theorem ([2] p. 484) let D denote the region

$$|t-t_0| \leq a, \|x - x_0\| \leq b, x = (x_1, \dots, x_n), x_0 = (x_{10}, \dots, x_{n0}) \tag{3.2}$$

and suppose that  $f(x, t)$  satisfies the Lipskhitz condition

$$\|f(x_1, t) - f(x_2, t)\| \leq k \|x_1 - x_2\| \tag{3.3}$$

whenever the pairs  $(x_1, t), (x_2, t)$  belong to D, where  $k$  is a positive constant. Then there is a constant  $\delta > 0$ , such that there exist a unique vector solution  $x(t)$  of the system (3.1) in the interval  $|t-t_0| \leq \delta$ . It is important to note ([2] p.485) that condition (3.3) is satisfied by requirement that  $\frac{\partial f_i}{\partial x_j}, i, j = 1, 2, \dots, n$  be continuous and bounded in D.

We now use the above theorem to prove the following theorem.

**Theorem 3.1**

Let  $n \geq 1, 0 < v_0 < \infty, \frac{\pi}{2} \leq \phi_0 \leq \frac{\pi}{2}$ . Then  $\exists$  a  $\delta > 0$  such that there is a unique vector solution  $(v, \phi)$  of the system (1.1) in the interval  $0 \leq t \leq \delta$ .

**Proof**

$$\frac{d}{dv} \left( -g \sin \phi + \frac{\sigma}{\rho} B_0^2 v - \frac{k}{m} v^n \right) = \frac{\sigma B_0^2}{\rho} - \frac{nk}{m} v^{n-1} \tag{3.4}$$

$$\frac{d}{d\phi} \left( -g \sin \phi + \frac{\sigma}{\rho} B_0^2 v - \frac{k}{m} v^n \right) = -g \cos \phi \tag{3.5}$$

also 
$$\frac{d}{dv} \left( -\frac{g \sin \phi}{v} \right) = \frac{g \cos \phi}{v^2}, \quad \frac{\partial}{\partial \phi} \left( \frac{-g \cos \phi}{v} \right) = \frac{g \sin \phi}{v} \tag{3.6}$$

Clearly (3.4) – (3.6) are continuous and bounded. Hence by the previous theorem, there exists a unique solution of the system.

**3.1 Exact Solutions**

Let  $L = \frac{v_0^2}{g}, T = \frac{v_0}{g}, v_0$  be units of length time, and velocity, then the non-dimensional equations are



$$\left. \begin{aligned} \frac{dv}{dt} &= -\sin \phi + \beta v - \alpha v^2, \quad v(0) = 1 \\ v \frac{d\phi}{dt} &= -\cos \phi, \quad \phi(0) = \phi_0 \\ \text{where } \alpha &= \sqrt{\frac{kv_0^2}{mg}}, \quad \beta = \frac{\alpha B_0^2 v_0}{g\rho} \end{aligned} \right\} \quad (3.7)$$

Consider vertical motion:  $\phi_0 = \frac{\pi}{2}$ .

In this special case, the particle is simply projected vertically upward. Let the time of ascent be  $t_a$ . Ascent phase  $0 < t < t_a$ . The initial and terminal conditions of this phase are  $v(0) = 1$ ,  $x = 0$ ,  $v(t_a) = 0$ , and  $\phi(0) = \frac{\pi}{2}$ ,  $y(0) = 0$ ,  $y(t_a) = h$ , then

$$\frac{dv}{dt} = -(1 - \beta v + \alpha^2 v^2), \quad v \frac{dv}{dy} = -(1 - \beta v + \alpha^2 v^2) \quad (3.8)$$

Thus, if (i)  $4\alpha^2 > \beta^2$ , then

$$v = \frac{B}{2\alpha^2} + \sqrt{\frac{4\alpha^2 - \beta^2}{2\alpha^2}} \left\{ \tan \left[ \frac{\sqrt{4\alpha^2 - \beta^2}(t_a - t)}{2} + \tan^{-1} \left( \frac{-\beta}{\sqrt{4\alpha^2 - \beta^2}} \right) \right] \right\}$$

(ii)  $4\alpha^2 = \beta^2$ , then  $v = \frac{1}{\alpha} - \left\{ \alpha^2 \left( t_a - t + \frac{1}{\alpha} \right) \right\}^{-1}$

(iii)  $4\alpha^2 < \beta^2$  then

$$\frac{\alpha^2}{\sqrt{\beta^2 - 4\alpha^2}} \left\{ \log_e \left| \frac{v - \frac{\beta - \sqrt{\beta^2 - 4\alpha^2}}{2\alpha^2}}{v - \frac{\beta + \sqrt{\beta^2 - 4\alpha^2}}{2\alpha^2}} \right| - \log_e \left| \frac{\frac{\beta + \sqrt{\beta^2 - 4\alpha^2}}{2\alpha^2}}{\frac{\beta - \sqrt{\beta^2 - 4\alpha^2}}{2\alpha^2}} \right| \right\} = \alpha^2 (t_a - t)$$

#### 4.0 Conclusion of the Results

When  $\beta = 0$ ,  $t_a = \frac{1}{\alpha} \tan^{-1}(\alpha)$  as obtained by Hayen [1]. When  $\alpha = 1$ ,  $t_a = 2$ , where as when  $\beta = 0$ ,  $\alpha = 1$ ,  $t_a = 1$ . When  $\beta^2 = 5\alpha^2$ ,  $\alpha = 1$ ,  $t_a = 0.98$ . Also, when  $\alpha = 1$ ,  $\beta = 2\alpha = 2$ , the particle cannot take off. So flight time increases only when  $4\alpha^2 > \beta^2$ . In other cases the electric field is a hindrance to the flight.

References

- [1] Hayen, J. C. (2003) Projectile motion in a resistant medium Part I: Exact solution and properties, Int. J. of Non-Linear Mechanics, 38: 357-369
- [2] Derrick, W. R. and Grossman, S. L. (1976) Differential Equations with Applications, Addison Wesley Reading.
- [3] Vajravelu, K. and Rivera, H. (2003) Hydromagnetic flow at an oscillating plate, Int. J. of Non-Linear Mechanics 38:305-312.

Consider vertical motion... In this special case the particle is simply projected vertically upward. Let the time of ascent be  $t_1$ . At the time  $t = t_1$ , the horizontal and vertical components of the velocity are  $v_x(t_1) = v_0 \cos \alpha$  and  $v_y(t_1) = 0$ .

(8.1) 
$$\left( \frac{v_0 \cos \alpha}{g} + t_1 - 1 \right) = \frac{v_0 \sin \alpha}{g} \left( \frac{v_0 \cos \alpha}{g} + t_1 - 1 \right) = \frac{v_0 \sin \alpha}{g} t_1$$

$$\left[ \frac{v_0 \cos \alpha}{g} + t_1 - 1 \right] = \frac{v_0 \sin \alpha}{g} t_1$$

- (ii)  $\alpha = \frac{\pi}{2}$ , then  $v = \frac{v_0}{2} - \frac{1}{2} g t^2$
- (iii)  $4\alpha^2 < \beta^2$  then

$$\left( \frac{v_0 \cos \alpha}{g} + t_1 - 1 \right) = \frac{v_0 \sin \alpha}{g} t_1$$

4.0 Conclusion of the results

When  $\beta = 0$ ,  $\alpha = \frac{\pi}{2}$  (i.e.  $\alpha = 90^\circ$ ) as obtained by Hayen [1]. When  $\alpha = \frac{\pi}{2}$ ,  $\beta = 2$ , where  $\beta$  is the ratio of the electric field to the magnetic field. Also, when  $\alpha = \frac{\pi}{2}$ ,  $\beta = 2\alpha = 2$ , the particle cannot take off. So the particle takes off only when  $\beta > \beta_c$ . In other cases the electric field is a hindrance to the flight.